

# ME8230

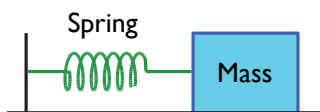
## Nonlinear Dynamics

Lecture I, part I  
Introduction, some basic math background,  
and some random examples

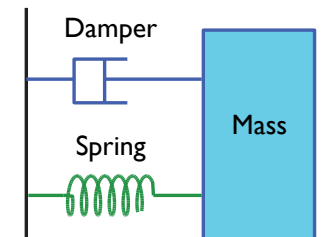
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## Spring mass damper system LINEAR

(unforced)



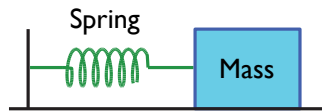
$$m\ddot{x} + kx = 0 \quad (\text{undamped})$$



$$m\ddot{x} + c\dot{x} + kx = 0 \quad (\text{damped})$$

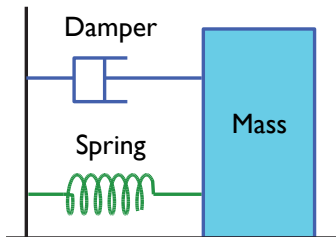
(You can add “external forcing terms” to all the “unforced” examples in this lecture )

# Spring mass damper system NON-LINEAR



undamped  $m\ddot{x} + f(x) = 0$

$f(x)$  = Nonlinear spring force  
 $g(\dot{x})$  = Nonlinear damping

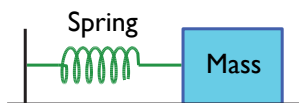


damped  $m\ddot{x} + g(\dot{x}) + f(x) = 0$

more general  
 nonlinear damped  
 mass-spring system  $m\ddot{x} + h(x, \dot{x}) = 0$

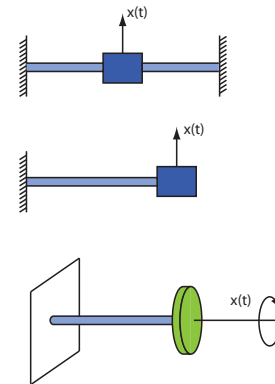
(unforced)

## Nonlinear spring-mass examples



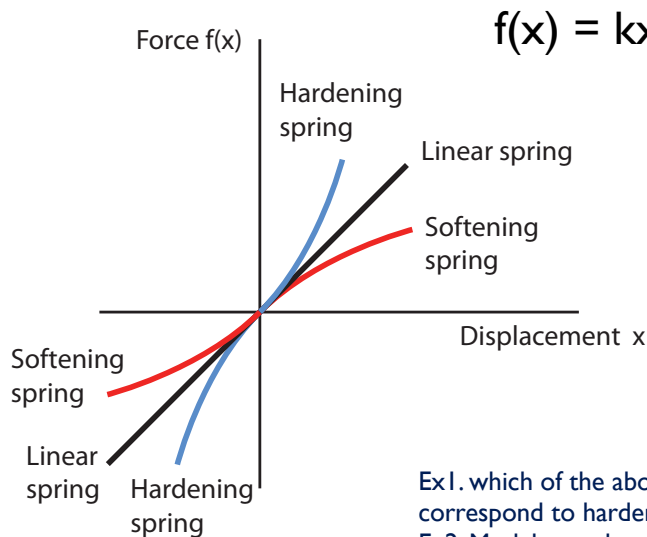
$$m\ddot{x} + f(x) = 0$$

can be a nonlinear  
 model of any of these  
 systems on the right  
when the displacements  
are large enough ... (but  
 mostly dominated by  
 the first mode)



Displacements small enough implies Linear spring approximation would be good enough (usually).

# Common Nonlinear Stiffness behaviors



$$f(x) = kx + ax^3, \quad a > 0, k > 0$$

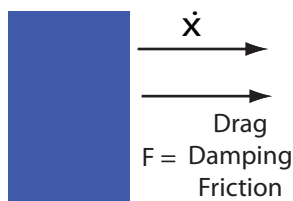
$$f(x) = kx$$

$$f(x) = kx - ax^3$$

Different nonlinearities might lead to different dynamical behaviors

Ex1. which of the above cubic  $f(x)$  correspond to hardening and softening springs?  
 Ex2. Model mass hanging from a taut horizontal string.

# Nonlinear damping examples



Linear damping:  $F = -c\dot{x}$

(negative sign indicates drag opposes velocity)

Fluid drag at large Re:  $F = -c\dot{x}|\dot{x}| = -c(\dot{x})^2 \text{sign}(\dot{x})$

Coulomb friction: "dry friction"  
 if  $\dot{x} = 0$ ,  $-\mu_s N \leq F \leq \mu_s N$   
 if  $\dot{x} \neq 0$ ,  $F = -\mu_d N \text{sign}(\dot{x})$

Material damping at high enough strain rates is nonlinear

# Many degrees of freedom systems

$$X = \begin{bmatrix} x_1 \\ x_2 \\ \cdot \\ \cdot \\ x_n \end{bmatrix}$$

## LINEAR

$$M\ddot{X} + C\dot{X} + KX = 0$$

Equations obtained from FEM, modal analysis, etc.

K - stiffness matrix

M - mass matrix

C - damping matrix

E.g., trusses/beam/plate/other elastic structures undergoing small deformations

## NON-LINEAR

$$M\ddot{X} + H(X, \dot{X}) = 0 \leftarrow$$

Pretty general: most smooth nonlinear (discrete, unforced) mechanical systems have such equations

# Forced finite dimensional mechanical systems

$$M(X) \ddot{X} = G(t, X, \dot{X})$$

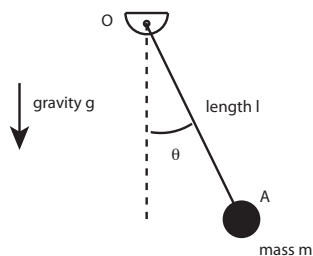
pretty general mechanical system

$$M(X) \ddot{X} = G_1(X, \dot{X}) + G_2(t)$$

special case  
e.g., double pendulum with joint torques

# Other Examples

## Simple pendulum



$$ml^2\ddot{\theta} + mgl \sin \theta = 0$$

$$\ddot{\theta} + \frac{g}{l} \sin \theta = 0$$

Do you know what the solutions  
to this equation looks like?  
Hint: Not sines and cosines in general.

# Aeroelastic oscillations

Tacoma Narrows Bridge Collapse



NASA Tail Flutter Test

Nonlinear (negative) damping

“Hopf bifurcation”

“Limit cycles”

Youtube videos of Tacoma Narrows Bridge Collapse and NASA tail flutter test

# Motions of disks and cylinders

Equations of motion for a cylinder rolling without slip

$$Q_{21}\ddot{\psi} + Q_{22}\ddot{\phi} + Q_{23}\ddot{\theta} = S_2, \quad i = 1, 2, 3. \quad (1)$$

where

$$Q_{11} = A \sin \phi - mHR \cos \phi + mH^2 \sin \phi,$$

$$Q_{12} = 0, \quad Q_{13} = -mHR,$$

$$Q_{21} = Q_{23} = 0, \quad Q_{22} = -mR^2 - mH^2 - A,$$

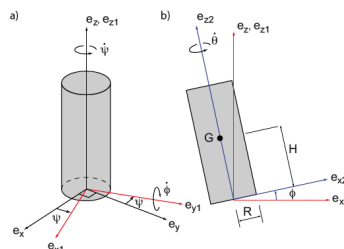
$$Q_{31} = C \cos \phi + mR^2 \cos \phi - mRH \sin \phi,$$

$$Q_{32} = 0, \quad Q_{33} = C + mR^2,$$

$$S_1 = (C - 2A - 2mH^2)\dot{\psi}\dot{\phi} \cos \phi + C\dot{\phi}\dot{\theta} - 2mHR\dot{\psi}\dot{\phi} \sin \phi,$$

$$S_2 = (C - A + mR^2 - mH^2)\dot{\psi}^2 \sin 2\phi/2 + (C + mR^2)\dot{\theta}^2 \sin \phi + mHR\dot{\psi}\dot{\theta} \cos \phi + mHR\dot{\psi}^2 \cos 2\phi + mg(R \cos \phi - H \sin \phi),$$

$$S_3 = C\dot{\psi}\dot{\phi} \sin \phi + 2mR\dot{\psi}\dot{\phi}(R \sin \phi + H \cos \phi).$$



Paper by Srinivasan and Ruina  
See demo, video and simulation.

# Double pendulum



(Double pendulum video: [Strogatz](#))

See also MATLAB  
example

(Hamiltonian) Chaos  
Sensitive dependence on initial  
conditions  
Lyapunov exponents

## Double pendulum with torque actuators at each joint and PID control

“Robot arm”

See MATLAB example

[Animations of controlled robot arm  
\(double/triple pendulum\)](#)

# Robots and humans

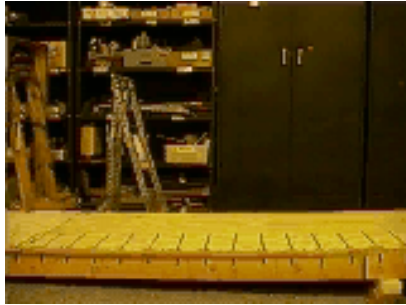
Hybrid (piecewise smooth) systems

Non-smooth systems

Limit cycles

Stability, Linearized

Collisions



Videos:

[Passive dynamic robots](#)

[Spring-mass hopper](#)

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ODEs corresponding to  
mechanical systems tend to be  
(most naturally)  
2nd order ODEs

Highest derivative order = 2



# High (e.g., 2nd) order ODEs to first order ODEs

always possible  
to convert

## Example

Say you are given 2nd order ODE:  $m\ddot{x} + kx = 0$

Introduce new variables  $x_1 = x$   
for all derivatives of  $x$ ,  $x_2 = \dot{x}$   
except the highest order (=2)

Write an ODE for each  $\dot{x}_1 = x_2$   
of the new variables  $\dot{x}_2 = -kx_1/m$

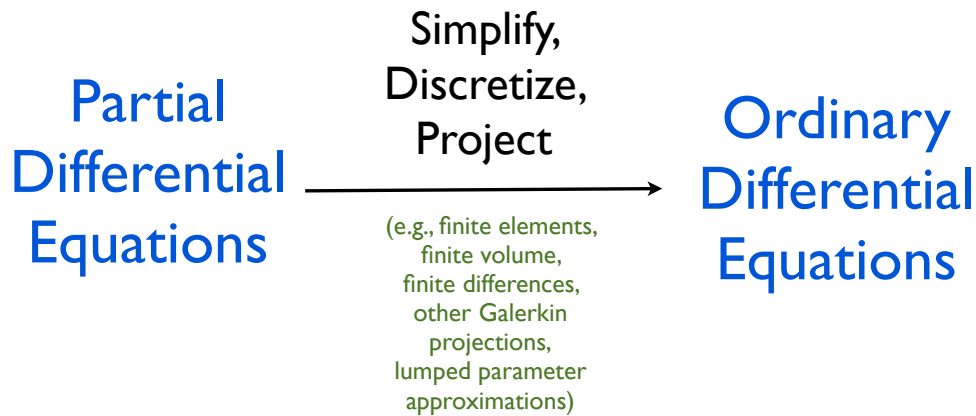
# What is this course about?

Mostly, ordinary  
differential equations  $\frac{dy}{dt} = f(y, t)$   
of the form

$t$  = independent variable  
 $y$  = dependent variable  
(can be scalar or vector)

$$y = \begin{bmatrix} y_1 \\ y_2 \\ \cdot \\ \cdot \\ \cdot \\ y_n \end{bmatrix}$$

# What to do with PDEs?



Some of the ideas in the course actually carry over to PDEs without reducing them to ODEs, but we will not discuss such.

## Linear ODEs

$$\frac{dy}{dt} = f(y, t)$$

Linear time-varying inhomogeneous (non-homogeneous) “variable coefficients”

$$\frac{dy}{dt} = A(t)y + B(t)$$

Inhomogeneous means  $B(t)$  is not zero

Linear time-varying homogeneous (variable coefficients)

$$\frac{dy}{dt} = A(t)y$$

Homogeneous means  $B(t)=0$

Linear time-invariant homogeneous (constant coefficients)

$$\frac{dy}{dt} = Ay$$

$A$  is a constant and not time-varying

## Linear superposition for linear homogeneous ODE

$$\frac{dy}{dt} = A(t)y$$

If  $y_1(t)$  and  $y_2(t)$  are solutions,  
then  $y_3 = y_1(t) + y_2(t)$  is also a solution.

$$\begin{array}{l} \frac{dy_1(t)}{dt} = A(t)y_1(t) \\ \frac{dy_2(t)}{dt} = A(t)y_2(t) \end{array} \xrightarrow{\text{add}} \frac{d(y_1(t) + y_2(t))}{dt} = A(t)y_1(t) + A(t)y_2(t) = A(t)(y_1(t) + y_2(t))$$

$$\text{Thus } \frac{y_3(t)}{dt} = A(t)y_3(t)$$

## Linear superposition for linear inhomogeneous ODE

$$\frac{dy}{dt} = A(t)y(t) + B(t)$$

**B(t)** - Input (the “inhomogeneous term”)

**y(t)** - Output

Say  $y_1(t)$  is a solution when  $B_1(t)$  is the input.

Say  $y_2(t)$  is a solution when  $B_2(t)$  is the input.

$y = \lambda_1 y_1 + \lambda_2 y_2$  is a solution when  $B = \lambda_1 B_1 + \lambda_2 B_2$ .

$$\begin{array}{l} \dot{y}_1 = A(t)y_1 + B_1(t) \\ \dot{y}_2 = A(t)y_2 + B_2(t) \end{array} \xrightarrow{\text{add}} \dot{y}_1 + \dot{y}_2 = A(t)(y_1 + y_2) + (B_1(t) + B_2(t))$$

# Key properties of linear systems

- Superposition of solutions
- proportionality (input and output scale together).
- Various analytical and semi-analytical techniques available for periodic, non-periodic forcing. In the absence of forcing, “closed-form” solutions known. e.g., Laplace transforms, Fourier series, Green’s functions, modal analysis, etc.

Nonlinear systems have none of these nice properties, in general. This course is about making sense of nonlinear ODEs by other means.